

In the present case we have to consider also the influence of pressure which was not considered in paper [2]. The relation (6) was derived under the assumption of a constant energy gap [5]. As mentioned above, we suppose a change of the energy gap W' or a change of the activation energy W caused by pressure. At the first approximation, in accordance with the relation (4) it follows

$$(8) \quad W(p) = W_0 - Cp.$$

Consequently in our case we can not make use of relation (6), but apply the more general

$$(9) \quad \sigma = \sigma_0 \exp [-W(p)/kT].$$

Developing the function $f(p, T) = W(p)/kT$ into the Taylor polynomial in the environment of the point (p_0, T_0) , and consider only the terms linearly dependent on temperature and pressure $p = \Delta p$, respectively, we receive with regard to (7) the expression

$$(10) \quad f(p, T) = \frac{W(p)}{kT} = \frac{W(p_0)}{kT_0} + \frac{1}{kT_0} \frac{\partial W(p)}{\partial p} p - \alpha(T - T_0).$$

As according to (4) $\partial W(p)/\partial p = -C$, hence from (9) and (10) follows the expression

$$(11) \quad \sigma = \sigma_{0p} \exp \alpha(T - T_0),$$

where

$$(12) \quad \sigma_{0p} = \sigma_{00} \exp \left[-\frac{W(p_0)}{kT_0} + \frac{C}{kT_0} \cdot p \right],$$

where σ_{0p} is the conductivity at temperature T_0 and pressure p and σ_{00} is the pre-exponential factor. In agreement with the relations (6) and (12) it will be

$$(13) \quad \sigma_{0p} = \sigma_0 \exp \left[\frac{C}{kT_0} p \right]$$

where σ_0 is the conductivity at $p = 0$ bar and $T = T_0$.

As it can be seen considering the pressure it becomes evident that relation (11) results instead of (6). Analogously with paper [2], taking in every case σ_{0p} instead of σ_0 , the following expression of V-A characteristic can be obtained

$$(14) \quad I = \frac{U}{R_{0p}} \exp [mU - nU^2],$$

where

$$(15) \quad m = Mad_i, \quad n = Nad_i, \quad d_i = D\sigma_{0p}.$$

The constants M, N, D are positive. The resistance R_{0p} occurring instead of R_0 from [2] is the resistance of the sample at the temperature $T = T_0$ and the pressure p .

Let us now investigate the dependence of the parameter m in the equation (14) on the pressure. It follows from the relations (7) and (8) that

$$(16) \quad \alpha = \frac{W_0 - Cp}{kT_0^2}.$$

Developing the exponential function in the relation (13), we shall obtain in the first approximation the expression

$$(17) \quad \sigma_{0p} = \sigma_0 \left(1 + \frac{C}{kT_0} \cdot p \right).$$

Using the relations (15), (16), (17) we shall arrive to the expression for m .

$$(18) \quad m = b(1 + k_1 p + k_2 p^2),$$

where

$$b = \frac{MD\sigma_0 W_0}{kT_0^2}, \quad k_1 = C \left(\frac{1}{kT_0} - \frac{1}{W_0} \right), \quad k_2 = -\frac{C^2}{kT_0 W_0}.$$

Assuming $nU^2 \ll mU$, [2] and taking into account the relation

$$(19) \quad \frac{1}{R_{0p}} = \frac{1}{R_{00}} (1 + Kp)$$

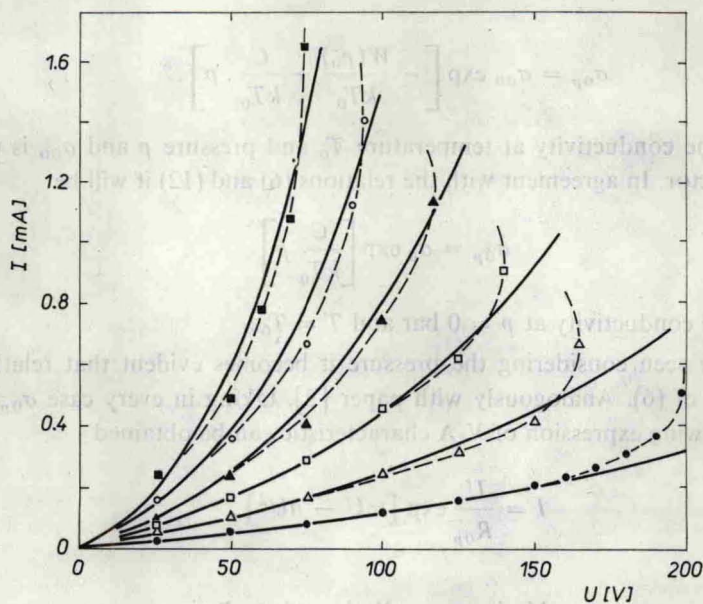


Fig. 1. The comparison of measured V-A characteristics with those expressed by relation (2) using suitable constants. ● — atmospheric pressure, △ — 1 kbar, □ — 2 kbar, ▲ — 3 kbar, ○ — 4 kbar, ■ — 5 kbar.